## Codes and Expansions (CodEx) Seminar

## Computing SIC-POVMs using

## Permutation Symmetries and Stark Units

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## Overview

- Zauner's conjecture
- numerical search
- symmetries \& Fibonacci-Lucas SIC-POVMs
- exact solutions from numerical solutions
- using overlaps
- using permutations
- numerical and exact solutions from Stark units
- conclusions \& outlook


## A Simple to State Problem

Are there $d^{2}$ vectors $\boldsymbol{v}_{1}, \boldsymbol{v}_{2}, \ldots, \boldsymbol{v}_{d^{2}} \in \mathbb{C}^{d}$ in the complex vector space of dimension $d$ such that:
$\begin{array}{rlrl}\text { (i) } & \left\langle\boldsymbol{v}_{j} \mid \boldsymbol{v}_{j}\right\rangle & =1 & \\ \text { for } j=1, \ldots, d^{2} \\ \text { (ii) } & \left|\left\langle\boldsymbol{v}_{j} \mid \boldsymbol{v}_{k}\right\rangle\right|^{2} & =\frac{1}{d+1} & \\ \text { for } 1 \leq j<k \leq d^{2}\end{array}$
The vectors $\boldsymbol{v}_{j}$ form an equiangular tight frame/finite unit norm tight frame.

## A Simple to State Problem

Are there $d^{2}$ vectors $\boldsymbol{v}_{1}, \boldsymbol{v}_{2}, \ldots, \boldsymbol{v}_{d^{2}} \in \mathbb{C}^{d}$ in the complex vector space of dimension $d$ such that:

$$
\begin{equation*}
\left\langle\boldsymbol{v}_{j} \mid \boldsymbol{v}_{j}\right\rangle=1 \quad \text { for } j=1, \ldots, d^{2} \tag{i}
\end{equation*}
$$

(ii) $\quad\left|\left\langle\boldsymbol{v}_{j} \mid \boldsymbol{v}_{k}\right\rangle\right|^{2}=\frac{1}{d+1} \quad$ for $1 \leq j<k \leq d^{2}$

The vectors $\boldsymbol{v}_{j}$ form an equiangular tight frame/finite unit norm tight frame.

All solutions form a real algebraic variety, using $2 d$ real variables per vector

$$
\boldsymbol{v}_{j}=\left(a_{j, 1}+i b_{j, 1}, a_{j, 2}+i b_{j, 2}, \ldots, a_{j, d}+i b_{j, d}\right)^{T} \quad(i=\sqrt{-1})
$$

$2 d^{3}$ variables, $d^{2}$ equations (i) of degree 2 and $\binom{d^{2}}{2}$ equations (ii) of degree 4 .

## Weyl-Heisenberg Group

- generators:

$$
H_{d}:=\langle X, Z\rangle
$$

$$
\text { where } X:=\sum_{j=0}^{d-1}|j+1\rangle\langle j| \text { and } Z:=\sum_{j=0}^{d-1} \omega_{d}^{j}|j\rangle\langle j|
$$

$$
\left(\omega_{d}:=\exp (2 \pi i / d)\right)
$$

- relations:

$$
\left(\omega_{d}^{c} X^{a} Z^{b}\right)\left(\omega_{d}^{c^{\prime}} X^{a^{\prime}} Z^{b^{\prime}}\right)=\omega_{d}^{a^{\prime} b-b^{\prime} a}\left(\omega_{d}^{c^{\prime}} X^{a^{\prime}} Z^{b^{\prime}}\right)\left(\omega_{d}^{c} X^{a} Z^{b}\right)
$$

- basis:

$$
H_{d} / \zeta\left(H_{d}\right)=\left\{X^{a} Z^{b}: a, b \in\{0, \ldots, d-1\}\right\} \cong \mathbb{Z}_{d} \times \mathbb{Z}_{d}
$$

trace-orthogonal basis of all $d \times d$ matrices

## Constructing SIC-POVMs

## Ansatz:

SIC-POVM that is the orbit under the Weyl-Heisenberg group $H_{d}$, i. e.,

$$
\begin{aligned}
\left|\boldsymbol{v}^{(a, b)}\right\rangle & :=X^{a} Z^{b}\left|\boldsymbol{v}^{(0,0)}\right\rangle \\
\left|\left\langle\boldsymbol{v}^{(a, b)} \mid \boldsymbol{v}^{\left(a^{\prime}, b^{\prime}\right)}\right\rangle\right|^{2} & = \begin{cases}1 & \text { for }(a, b)=\left(a^{\prime}, b^{\prime}\right) \\
1 /(d+1) & \text { for }(a, b) \neq\left(a^{\prime}, b^{\prime}\right)\end{cases} \\
\left|\boldsymbol{v}^{(0,0)}\right\rangle & =\sum_{j=0}^{d-1}\left(x_{2 j}+i x_{2 j+1}\right)|j\rangle, \\
\left(x_{0}, \ldots, x_{2 d-1}\right. \text { are real variables, } & \left.x_{1}=0\right)
\end{aligned}
$$

$\Longrightarrow$ we have to find only one fiducial vector $\left|\boldsymbol{v}^{(0,0)}\right\rangle$ instead of $d^{2}$ vectors
$\Longrightarrow$ polynomial equations with $2 d-1$ variables, but already quite complicated for $d=6$

## Jacobi Group (or Clifford Group)

- automorphism group of the Weyl-Heisenberg group $H_{d}$, i. e.

$$
\forall T \in J_{d}: T^{\dagger} H_{d} T=H_{d}
$$

- the action of $J_{d}$ on $H_{d}$ modulo phases corresponds to the symplectic group $\mathrm{SL}\left(2, \mathbb{Z}_{d}\right)$, i.e.

$$
T^{\dagger} X^{a} Z^{b} T=\omega_{d}^{c} X^{a^{\prime}} Z^{b^{\prime}} \quad \text { where }\binom{a^{\prime}}{b^{\prime}}=\tilde{T}\binom{a}{b}, \tilde{T} \in \mathrm{SL}\left(2, \mathbb{Z}_{d}\right)
$$

$\Longrightarrow$ homomorphism $J_{d} \rightarrow \mathrm{SL}\left(2, \mathbb{Z}_{d}\right)$

- additionally: complex conjugation (anti-unitary)

$$
X^{a} Z^{b} \mapsto X^{a} Z^{-b} \quad \text { corresponding to }\left(\begin{array}{cc}
1 & 0 \\
0 & -1
\end{array}\right)
$$

## Zauner's Conjecture

[G. Zauner, Dissertation, Universität Wien, 1999]

## Conjecture:

For every dimension $d \geq 2$ there exists a SIC-POVM whose elements are the orbit of a rank-one operator $E_{0}$ under the Weyl-Heisenberg group $H_{d}$. What is more, $E_{0}$ commutes with an element $S$ of the Jacobi group $J_{d}$. The action of $S$ on $H_{d}$ modulo the center has order three.
support for this conjecture (to date):

- numerical solutions for all dimensions $d \leq 193$, plus a few more
- exact algebraic solutions for some dimensions (see below)
one of the prize problems in
[Paweł Horodecki, Łukasz Rudnicki, Karol Życzkowski, Five open problems in quantum information, arXiv:2002.03233]


## Numerical Search for SIC-POVMs

- "second frame potential" for 2-designs

$$
\left.\sum_{i, j=1}^{d^{2}}\left|\left\langle v^{(i)} \otimes v^{(i)} \mid v^{(j)} \otimes v^{(j)}\right\rangle\right|^{2}=\sum_{i, j=1}^{d^{2}}\left|\left\langle v^{(i)} \mid v^{(j)}\right\rangle\right|^{4}=d^{2} \sum_{a, b=1}^{d}\left|\langle\psi| X^{a} Z^{b}\right| \psi\right\rangle\left.\right|^{4}
$$

- for any state $|\psi\rangle \in \mathbb{C}^{d}$

$$
\begin{array}{rlr}
f(|\psi\rangle) & =\sum_{j, k=1}^{d}\left|\sum_{\ell=1}^{d}\langle\psi \mid j+\ell\rangle\langle\ell \mid \psi\rangle\langle\psi \mid k+\ell\rangle\langle j+k+\ell \mid \psi\rangle\right|^{2} & \\
& =\sum_{j, k=1}^{d}|\underbrace{\sum_{\ell=1}^{d} \bar{\psi}_{j+\ell} \psi_{\ell} \bar{\psi}_{k+\ell} \psi_{j+k+\ell}}_{=: G(j, k)}|^{2} & \geq \frac{2}{d+1}
\end{array}
$$

with equality iff $|\psi\rangle$ is a fiducial vector for a Weyl-Heisenberg SIC-POVM

- gradient descent to minimize $f(|\psi\rangle)$, subject to unit norm


## Numerical Search for SIC-POVMs

- for any state $|\psi\rangle \in \mathbb{C}^{d}$

$$
f(|\psi\rangle)=\sum_{j, k=1}^{d}\left|\sum_{\ell=1}^{d} \bar{\psi}_{j+\ell} \psi_{\ell} \bar{\psi}_{k+\ell} \psi_{j+k+\ell}\right|^{2} \geq \frac{2}{d+1}
$$

with equality iff $|\psi\rangle$ is a fiducial vector for a Weyl-Heisenberg SIC-POVM

- gradient descent to minimize $f(|\psi\rangle)$, subject to unit norm
- use $F(\vec{x})=f\left(\frac{P \vec{x}}{\|P \vec{x}\|}\right)$ for an arbitrary vector $\vec{x} \in \mathbb{C}^{d}$, where $P$ is the projection onto a subspace (prescribed symmetry)
- chain rule yields a relatively simple formula for the gradient of $F(\vec{x})$ in terms of the gradient of $f$
- complexity $\mathcal{O}\left(d^{3}\right)$ for both the function and the gradient when storing $\mathcal{O}\left(d^{2}\right)$ intermediate values


## Numerical Search for SIC-POVMs

- efficient implementation of $F(\vec{x})$ and its gradient in C++ by Andrew Scott
- parallel computation of the function/gradient using OpenMP/CUDA
- minimization using limited-memory Broyden-Fletcher-Goldfarb-Shanno (BFGS) algorithm
- search runs into local minima, we need many random initial points
- running many instances on HPC clusters by MPG and GWDG
-     - for $d=189$ : approx. $23.3 \times 10^{6}$ trials, 3.48 CPU years
- for $d=190$ : approx. $66.8 \times 10^{6}$ trials, 10.51 CPU years
- for $d=193$ : approx. $78.3 \times 10^{6}$ trials, 13.00 CPU years
- for $d=5779: 55065$ trials, 17.69 GPU years, no success


## Average Number of Iterations



## Average Number of Iterations with Zauner Symmetry



## Fibonacci-Lucas SIC-POVMs

[Markus Grassl \& Andrew J. Scott, JMP 58, December 2017, arXiv:1707.02944]

- (exact) symmetry analysis of a numerical solution for $d=124$ $\Longrightarrow$ symmetry group of order 30 (prescribed order 6)
- identified as part of a series of dimensions (related to Lucas numbers) $d=4,8,19,48,124,323,844,2208,5779,15128,39604, \ldots$
- symmetry group of order $6 k$ related to Fibonacci numbers, $F=\left(\begin{array}{ll}0 & 1 \\ 1 & 1\end{array}\right)$
- new exact solutions for $d=124$ and $d=323$ (previously $d=48$ ) (found using the symmetries and via Gröbner bases)
- new numerical solutions for $d=844$ and $d=2208$ (previously $d=323$ )
- generalisations related to generalised Fibonacci/Lucas numbers, using $A_{m}=\left(\begin{array}{ll}0 & 1 \\ 1 & m\end{array}\right)$


## Symmetries of SIC-POVMs



## SIC-POVMs by Numerical Search



## Ray Class Field Conjecture

[Appleby, Flammia, McConnell \& Yard, arXiv:1604.06098 \& arXiv:1701.052000]
CodEx talks by M. Appleby, S. Flammia, G. Kopp

## Ray class field conjecture

let $\mathbb{E}$ be the field containing all rank-one projection operators of a SIC-POVM

$$
\mathbb{Q} \triangleleft \mathbb{K}=\mathbb{Q}(\sqrt{D}) \triangleleft \mathbb{E}_{0} \triangleleft \mathbb{E}_{1} \triangleleft \mathbb{E}
$$

for the minimal field:

- $\mathbb{E}$ is the ray class field over $\mathbb{Q}(\sqrt{D})$ with conductor ${ }^{\text {a }} d^{\prime}$ with ramification at both infinite places, $D$ is the squarefree part of $(d+1)(d-3)$
- $\mathbb{E}_{1}$ contains the overlap phases and equals the ray class field with ramification only allowed at the infinite place taking $\sqrt{D}$ to a positive real number
- $\mathbb{E}_{0}$ is the Hilbert class field $H_{\mathbb{K}}$, in particular $h=\left[\mathbb{E}_{0}: \mathbb{K}\right]$ equals the class number of $\mathbb{K}$

$$
{ }^{\mathrm{a}} d^{\prime}=d, \text { or } d^{\prime}=2 d \text { for } d \text { even }
$$

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$$

- "Fact 8": $\operatorname{Gal}\left(\mathbb{E}_{1} / \mathbb{E}_{0}\right)$ permutes the overlaps.

For each $\sigma \in \operatorname{Gal}\left(\mathbb{E}_{1} / \mathbb{E}_{0}\right)$ there is a matrix $G_{\sigma} \in \mathrm{GL}\left(2, \mathbb{Z} / d^{\prime} \mathbb{Z}\right)$ such that ${ }^{\text {a }}$

$$
\sigma\left(\langle\psi| D_{\boldsymbol{p}}|\psi\rangle\right)=\langle\psi| D_{G_{\sigma} \boldsymbol{p}}|\psi\rangle
$$

$G_{\sigma}$ commutes with matrices $F$ related to symmetries $U_{F}$ of the fiducial vector $|\psi\rangle$.

$$
{ }^{\mathrm{a}} D_{\boldsymbol{p}}=D_{a, b}=\left(e^{\frac{i \pi}{d}}\right)^{a b} X^{a} Z^{b}
$$

## Exact Solutions from Numerical Solutions

[Appleby, Chien, Flammia \& Waldron, J. Phys. A. 51, 2018, arXiv:1703.05981]

- matrix group $\mathcal{M}=\left\{G_{\sigma}: \sigma \in \operatorname{Gal}\left(\mathbb{E}_{1} / \mathbb{E}_{0}\right)\right\}$, commutes with the symmetry
- projection operator $\Pi=|\psi\rangle\langle\psi|$

$$
\text { "Fact 8:" } \quad \sigma\left(\operatorname{Tr}\left(\Pi D_{\boldsymbol{p}}\right)\right)=\operatorname{Tr}\left(\Pi D_{G_{\sigma} \boldsymbol{p}}\right)
$$

- expansion coefficients $c_{\boldsymbol{p}}=\operatorname{Tr}\left(\Pi D_{\boldsymbol{p}}\right)$ in the same orbit under $\mathcal{M}$ are related by Galois conjugation
- the coefficients of the polynomial $f_{\boldsymbol{p}_{0}}(z)=\prod_{\boldsymbol{p} \in \boldsymbol{p}_{0}^{M}}\left(z-c_{\boldsymbol{p}}\right)$ lie in a number field of "small" degree
- find the exact minimal polynomials of those coefficients (requires high-precision numerical solution)
- find the roots of the exact polynomials $f_{p_{0}}(z)$ in the ray class field
- compute $\Pi$ from the $d^{2}$ expansion coefficients $c_{\boldsymbol{p}}$
- exact solutions for some $d \leq 48$ ( $d \leq 100$ work in progress)


## More Exact Solutions from Numerical Solutions

[Markus Grassl, Exact SIC-POVMs from permutation symmetries, in preparation]

- when $G_{\sigma}$ has determinant 1 , there exists a unitary $U_{G_{\sigma}}:=T_{\sigma}$ with

$$
\sigma\left(\operatorname{Tr}\left(\Pi D_{\boldsymbol{p}}\right)\right)=\operatorname{Tr}\left(\Pi D_{G_{\sigma} \boldsymbol{p}}\right)=\operatorname{Tr}\left(\Pi T_{\sigma} D_{\boldsymbol{p}} T_{\sigma}^{\dagger}\right)=\operatorname{Tr}\left(T_{\sigma}^{\dagger} \Pi T_{\sigma} D_{\boldsymbol{p}}\right)
$$

$\Longrightarrow$ action of $T_{\sigma}^{\dagger}$ on the projection $\Pi$ and on the state $|\psi\rangle$

- when $G_{\sigma}=\left(\begin{array}{cc}\alpha & 0 \\ 0 & \alpha^{-1}\end{array}\right)$ is additionally diagonal, then $T_{\sigma}$ is a permutation matrix
- moreover, assume that $\sigma\left(D_{\boldsymbol{p}}\right)=D_{\boldsymbol{p}}$; then

$$
\sigma(\Pi)=T_{\sigma}^{\dagger} \Pi T_{\sigma}
$$

and hence

$$
\sigma\left(\Pi_{j, k}\right)=\Pi_{\alpha j, \alpha k}
$$

where the indices are computed modulo $d$

## More Exact Solutions from Numerical Solutions

- for the first column of $\Pi$ we have

$$
\sigma\left(\Pi_{j, 0}\right)=\Pi_{\alpha j, 0} \quad \text { for } j=0, \ldots, d-1
$$

- we can take the first column as (unnormalised) fiducial vector $\boldsymbol{v}$, unless it is zero (which was observed for $d=26,28,62,98,228$ ) $\Longrightarrow \sigma$ permutes the components of the fiducial vector, stabilising the first coordinate
- when the first column is zero, consider a non-zero column $k$ :

$$
\sigma\left(\Pi_{j, k}\right)=\Pi_{\alpha j, \alpha k} \stackrel{(*)}{=} \gamma \Pi_{\alpha j, k} \quad \text { for } j=0, \ldots, d-1
$$

$\Longrightarrow \sigma$ gives rise to a projective permutation action
$\Longrightarrow$ consider the action on ratios $v_{j} / v_{j^{\prime}}$
${ }^{(*)} \Pi$ has rank one, so column $\alpha k$ is proportional to column $k$, i.e., $\Pi_{j, \alpha k}=\gamma \Pi_{j, k}$

## More Exact Solutions from Numerical Solutions

outline of the procedure:

- compute a numerical fiducial vector with prescribed symmetry $S$
- determine the diagonal matrices $G_{\sigma} \in \mathrm{SL}\left(2, \mathbb{Z} / d^{\prime} \mathbb{Z}\right)$ in the centraliser of $S$
- the diagonal matrices correspond to a subgroup $H \leq\left(Z / d^{\prime} \mathbb{Z}\right)^{\times}$
- consider the rescaled fiducial vector ${ }^{\text {a }} \boldsymbol{v}$ with $v_{0}=1$
- the coefficients of the polynomial $f_{j}(z)=\prod_{\alpha \in H}\left(z-v_{\alpha j}\right)$ lie in a number field of "small" degree, fixed by (a subgroup of) the Galois group
- similar as before, find the exact coefficients of $f_{j}(z)$ from a high-precision numerical solution, and then compute its exact roots
$\Longrightarrow$ only $\mathcal{O}(d)$ numbers in a field of smaller degree

[^0]
## More Exact Solutions from Numerical Solutions

- the assumption that $\sigma\left(D_{\boldsymbol{p}}\right)=D_{\boldsymbol{p}}$ appears to be true
- new exact solutions for 57 additional dimensions (so far)

$$
\begin{aligned}
d= & 26,38,42,49,52,57,61,62,63,65,67,73,74,78,79,84,86,91,93,95, \\
& 97,98,103,109,111,122,127,129,133,134,139,143,146,147 \\
& 151,155,157,163,168,169,172,181,182,183,193,199 \\
& 201,228,259,292,327,364,399,403,489,844,1299
\end{aligned}
$$

- fiducial vectors lie in a proper ("small") subfield of the ray class field from before, that intersects with the cyclotomic field $\mathbb{Q}\left(\zeta_{d^{\prime}}\right)$ trivially or in a smaller cyclotomic field
- "small ray class field conjecture":

The minimal field containing a (suitably rescaled) fiducial vector is a ray class field whose conductor is a particular factor of the ideal $d \mathcal{O}_{\mathrm{IK}}$ with ramification at one of the infinite places.

## Prime Dimensions $p \equiv 1 \bmod 3$

- for prime dimensions $d=p \equiv 1 \bmod 3$, the Zauner symmetry $F_{z}$ is conjugate to a diagonal matrix $\widetilde{F}_{z}$
- the centraliser of $\widetilde{F}_{z}$ contains all diagonal matrices in $\operatorname{SL}(2, \mathbb{Z} / d \mathbb{Z})$
- the components $v_{j}, j=1, \ldots, d-1$, of the fiducial vector (with $v_{0}=1$ ) are on a single orbit with respect to the Galois group, i.e.,

$$
v_{\theta^{k}}=\sigma^{k}\left(v_{1}\right)
$$

for generators $\theta$ and $\sigma$ of $(\mathbb{Z} / d \mathbb{Z})^{\times}$and the Galois group, resp.

- for a permutation symmetry of order $3 \ell$, we need only $m=\frac{d-1}{3 \ell}$ numbers


## dream:

find a direct way to determine the algebraic number $v_{1}$, as well as $\sigma$ and $\theta$

## Prime Dimensions $p=n^{2}+3$

[Appleby, Bengtsson, Grassl, Harrison, McConnell, "SIC-POVMs from Stark Units"]

## Conjecture:

- for prime dimensions $p=n^{2}+3(n>0)$, there is an almost flat fiducial vector $\boldsymbol{v}$ with

$$
v_{j}= \begin{cases}-2-\sqrt{d+1} & j=0 \\ \sqrt{v_{0} e^{i \vartheta_{j}}} & j>0\end{cases}
$$

- the components of $\boldsymbol{v}$ generate a "small" ray class field $\mathbb{K}^{\mathfrak{m}}$ with finite modulus $\sqrt{d+1} \pm 1$ and ramification at one infinite place
- the phases $e^{i \vartheta_{j}}$ are Galois conjugates of (real) Stark units for the ray class field $\mathbb{K}^{\mathfrak{m}}$


## Application of Stark's Conjectures

- for certain ray class fields $\mathbb{K}^{\mathfrak{m}}$ over the real quadratic field $\mathbb{K}=\mathbb{Q}(\sqrt{D})$, $D>0$, one can compute numerical approximations of Stark units $\epsilon_{\sigma}$ via special values of derivatives of $L$-functions
- the Stark units are labelled by elements $\sigma$ of the Galois group $\operatorname{Gal}\left(\mathbb{K}^{\mathfrak{m}} / \mathbb{K}\right)$ such that $\epsilon_{\sigma}=\sigma\left(\epsilon_{0}\right)$
- from numerical Stark units with sufficiently high precision, we can deduce their exact minimal polynomial over $\mathbb{K}$
- we have a heuristic that allows us to deduce the required precision from numerical Stark units with low precision
- the complexity of the calculation appears to be roughly $\mathcal{O}\left(\operatorname{deg}\left(\mathbb{K}^{\mathfrak{m}} / \mathbb{K}\right) \times(\# \text { digits })^{3.3}\right)$


## Runtime $L$-Functions

total CPU time to compute the numerical derivative of $L$-functions using Magma and PARI/GP (last three cases)

| $d$ | $\operatorname{deg}\left(\mathbb{K}^{\mathfrak{m}} / \mathbb{K}\right)$ | log height | precision | CPU time |
| :---: | :---: | ---: | :---: | :---: |
| 487 | 324 | 424 | 1000 digits | 251 hours |
| 787 | 262 | 299 | 1000 digits | 118 hours |
| 2707 | 902 | 1861 | 3800 digits | 900 days |
| 4099 | 1366 | 974 | 2000 digits | 170 days |
| 5779 | 214 | 127 | 300 digits | 18 min |
| 1447 | 964 | 2158 | 4600 digits | 111 days |
| 19603 | 2178 | 1754 | 4000 digits | 82 days |
| 2503 | 3336 | 6464 | 13000 digits | 60.5 years |

## Flipping the Sign

- real quadratic field $\mathbb{K}=\mathbb{Q}(\sqrt{D})$ with non-trivial automorphism $\tau: \sqrt{D} \mapsto-\sqrt{D}$
- embedding $\mathfrak{j}: \mathbb{K} \hookrightarrow \mathbb{R}, \mathfrak{j}(\sqrt{D})>0, \mathfrak{j}^{\tau}(\sqrt{D})=\mathfrak{j}\left((\sqrt{D})^{\tau}\right)<0$
- "real" Stark units $\epsilon_{\sigma}: \mathfrak{j}\left(\epsilon_{\sigma}\right)>0$
- "complex" Stark units $\epsilon_{\sigma}^{\tau}: \mathfrak{j}\left(\epsilon_{\sigma}^{\tau}\right)=\mathfrak{j}^{\tau}\left(\epsilon_{\sigma}\right) \in \mathbb{C} \backslash \mathbb{R}$
- minimal polynomial of $\epsilon_{\sigma}: \quad p_{1}(t) \in \mathbb{K}[t]$
$\Longrightarrow$ minimal polynomial of $\epsilon_{\sigma}^{\tau}: p_{2}(t)=p_{1}^{\tau}(t)$
- obstacle:
operation of $\sigma$ on $\epsilon_{0}^{\tau}$ would require factoring $p_{2}(t)$

for simplicity, we assume in the following class number $h=1$, i.e., $H_{\mathbb{K}}=\mathbb{K}$


## The Galois Polynomial

- fixing some labelling, we know how $\sigma$ permutes the $m$ (numerical) Stark units $\epsilon_{j}: \sigma\left(\epsilon_{j}\right)=\sigma_{\pi_{\sigma}(j)}$ for some permutation $\pi_{\sigma}$
- there exists a unique polynomial $g_{1}$ of degree at most $m-1$ such that

$$
\begin{equation*}
g_{1}\left(\epsilon_{j}\right)=\epsilon_{\pi_{\sigma}(j)} \quad \text { for } j=1, \ldots, m \tag{1}
\end{equation*}
$$

- using Newton interpolation, $g_{1}$ can be computed with $\mathcal{O}\left(m^{2}\right)$ arithmetic operations $\left(\mathcal{O}\left(m(\log m)^{2}\right)\right.$ when using FFT-based methods)
- the coefficients of $g_{1}$ are in $\mathbb{K}$, as $(1)$ is invariant wrt. $\operatorname{Gal}\left(\mathbb{K}^{\mathfrak{p} j^{\tau}} / \mathbb{K}\right)$
- $g_{2}(t)=g_{1}^{\tau}(t)$ corresponds to the action of $\sigma$ on $\epsilon_{j}^{\tau}: g_{2}\left(\epsilon_{j}^{\tau}\right)=\epsilon_{\pi_{\sigma}(j)}^{\tau}$
- potential computational obstacle:
we don't know an a priori bound for the required precision (for $d=19603$, the coefficients have more than 1 million digits)


## Solving the Sign Problem

Recall: We conjecture that the components of the fiducial vector are square roots of Galois conjugates of Stark units, i.e., $v_{\theta^{k}}=\sqrt{v_{0} \sigma^{k}\left(\epsilon_{0}^{\tau}\right)}$.
Problem: there are two square roots $\pm \sqrt{v_{0} \sigma^{k}\left(\epsilon_{0}^{\tau}\right)}$

## Solution:

- it turns out that polynomial $p_{2}\left(t^{2} / v_{0}\right)$ with $v_{0}=-2-\sqrt{d+1}$ factors in $\mathbb{K}[t]$ as

$$
v_{0}^{m} p_{2}\left(t^{2} / v_{0}\right)=p_{4}(t) p_{4}(-t)
$$

- pick the factor $p_{4}(t)$ and check which of the square roots is a root of $p_{4}(t)$
- we are left with a global sign ambiguity, i.e., two possibilities
- note: it does not matter which of the Galois conjugates of the Stark units is assigned to $\epsilon_{0}^{\tau}$; all choices yield eventually fiducial vectors


## Final Step: Combinatorial Search

so far, we have

- exact minimal polynomials $p_{1}(t), p_{2}(t), p_{4}(t) \in \mathbb{K}[t]$ and exact Galois polynomials $g_{1}(t), g_{2}(t) \in \mathbb{K}[t]$
- numerical square roots $\sqrt{v_{0} \epsilon_{j}^{\tau}}$ (up to a global sign) together with the permutation action of the (cyclic) Galois group $\operatorname{Gal}\left(\mathbb{K}^{\mathfrak{m}} / H_{\mathrm{K}}\right)$ on them final step:
- we have to identify which primitive element $\theta \in(\mathbb{Z} / d \mathbb{Z})^{\times}$corresponds to the action of $\sigma$
- we have to fix the global sign (we can choose the sign of the first coordinate)
- compute a (numerical) vector $\boldsymbol{v}$ for all choices (less than $d$ ) and test the overlap $|\langle\boldsymbol{v}| X| \boldsymbol{v}\rangle\left.\right|^{2} /\|\boldsymbol{v}\|^{4} \stackrel{?}{=} \frac{1}{d+1}$
- we know that $\sigma^{m / 2}$ corresponds to complex conjugation


## Exact Solution

We can also compute an exact representation of the fiducial vector without explicit factorisation in the extension field:

- define the field $\mathbb{L}=H_{\mathbb{K}}(\gamma)$ with $p_{4}(\gamma)=0$
- compute the exact Galois polynomial $g_{4}(t) \in H_{\mathrm{K}}[t]$ from the numerical values $\sqrt{v_{0} \epsilon_{\sigma}^{\tau}}$
- the action of the Galois automorphism $\sigma$ on $\mathbb{L}$ is defined by $\sigma: \gamma \mapsto g_{4}(\gamma)$
- we can compute the components of the fiducial vector using

$$
v_{0}= \pm(2+\sqrt{d+1}), \quad v_{1}=\gamma, \quad \text { and } v_{\theta j}=g_{4}\left(v_{j}\right) \text { for } j>0
$$

computational obstacles: missing an a priori bound on the precision to compute the exact Galois polynomial $p_{4}(t)$ and arithmetic in the field $\mathbb{L}$ is slow when the degree is large (use tower of subfields if possible)

## Verification of the Solution

- second frame potential for a fiducial vector

$$
f(|\psi\rangle)=\sum_{j, k=1}^{d}|\underbrace{\sum_{\ell=1}^{d} \bar{\psi}_{j+\ell} \psi_{\ell} \bar{\psi}_{k+\ell} \psi_{j+k+\ell}}_{=: G(j, k)}|^{2}=\frac{2}{d+1}
$$

- moreover $G(j, k)=\frac{\delta_{j, 0}+\delta_{k, 0}}{d+1}$
- $G(j, k)$ has an eightfold symmetry
- we don't need $d$-th roots of unity
- $\mathcal{O}\left(d^{3}\right)$ arithmetic operations
verifying the solution takes longer than computing it


## Runtime Verification

CPU time for the exact/numerical verification of the solution

| $d$ | $\operatorname{deg}\left(\mathbb{K}^{\mathfrak{m}} / \mathbb{K}\right)$ | precision | CPU time | $G(j, k)$ |
| ---: | ---: | ---: | :---: | :---: |
| 103 | $2^{2} \times 17$ | exact | 440 s | 1.3 s |
| 199 | $2 \times 11$ | exact | 310 s | 0.3 s |
| 487 | $2^{2} \times 3^{4}$ | exact | 31 days | 315 s |
| 787 | $2 \times 131$ | 10000 digits | 3 hours | 65 min |
| 1447 | $2^{2} \times 241$ | 10000 digits | 17.0 hours |  |
| 2707 | $2 \times 11 \times 41$ | 2000 digits | 11.2 hours |  |
| 4099 | $2 \times 683$ | 2000 digits | 36.5 hours |  |
| 5779 | $2 \times 107$ | 2000 digits | 100 hours | 88 min |
| 19603 | $2 \times 3^{2} \times 11^{2}$ | 1000 digits | 1367 days |  |
| 39604 | $2^{2} \times 3^{2} \times 5^{2}$ | 100 digits | 684 days | $\approx 28$ days |

## Solutions for $d=n^{2}+3$

- the method can be generalised to composite dimensions $d=n^{2}+3$
- even dimensions $d=n^{2}+3$ are divisible by 4 , but not by 8 ; almost flat fiducial vector after change of basis
- for composite dimensions, one has to compute Stark units for certain subfields as well
- there are more possibilities to match the action of $(\mathbb{Z} / d \mathbb{Z})^{\times}$and the action of the Galois group
so far, our method has been successfully applied in 34 dimensions:
$d=7,12,19,28,39,52,67,84,103,124,147,172,199,259,292,327,403,487$, $628,787,844,964,1027,1228,1299,1447,1684,1852,2404,2707,4099,5779$, 19603, and 39604


## Conclusions \& Outlook

- deterministic procedure to compute SIC-POVMs from Stark units
- successfully applied in 34 dimensions $d=n^{2}+3$; did not fail in any
- can we obtain a fiducial vector directly from the real Stark units, without "flipping the sign"?
- can we work with lower precision?
- can we avoid the combinatorial search in the final step?
- assuming Stark's conjectures to be true, can be prove that our construction always works?
- can we extend the method to other dimensions?
forthcoming publication:
Marcus Appleby, Ingemar Bengtsson, Markus Grassl, Michael Harrison, Gary McConnell, "SIC-POVMs from Stark Units"


## Thank you! Danke! Merci! Dziekuje!

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[^0]:    ${ }^{\text {a assuming }} v_{0} \neq 0$ for simplicity here

