

Math 670: Dg 12

Def: The ~~exterior~~ **exterior derivative** $d: \Omega^k(M) \rightarrow \Omega^{k+1}(M)$ is the unique anti-derivation of degree +1 s.t.

① $\forall f \in \Omega^0(M) = C^0(M)$, df is the differential of f .

② $d \circ d = 0$.

Prop: There is such a best.

Proof: In local coords, $\omega_p = \sum f_I dx_I$, and we define

$$(d\omega)_p = \sum (df_I) \wedge dx_I$$

We saw last time why this is a reasonable way to define the ~~exterior~~ derivative.

Independence of coordinates will follow from uniqueness, so it remains to show:

(i) This is an anti-derivation

(ii) It satisfies the cycle condition ②

(iii) uniqueness.

(i) & (ii) are both computations. Here's the computation for (ii):

If $\omega_p = \sum f_I dx_I$, then

$$(d\omega)_p = \sum_I df_I \wedge dx_I = \sum_I \left(\sum_i \frac{\partial f}{\partial x_i} dx_i \right) \wedge dx_I,$$

so

$$(dd\omega)_p = \sum_I \left(\sum_i d\left(\frac{\partial f}{\partial x_i}\right) \wedge dx_i \right) \wedge dx_I = \sum_I \left(\sum_i \left(\sum_j \frac{\partial^2 f}{\partial x_j \partial x_i} dx_j \right) \wedge dx_i \right) \wedge dx_I = 0$$

since $\frac{\partial^2 f}{\partial x_j \partial x_i} = \frac{\partial^2 f}{\partial x_i \partial x_j}$ but $dx_j \wedge dx_i = -dx_i \wedge dx_j$.

For (iii), suppose $\exists$ a degree 1 anti-derivation $D: \Omega^k(M) \rightarrow \Omega^{k+1}(M)$ satisfying ① & ②. In particular, $Df = df \forall f \in C^0(M)$.

Then $D(dx_I) = D(dx_{i_1} \wedge \dots \wedge dx_{i_k}) = \sum_j (-1)^{j-1} dx_{i_1} \wedge \dots \wedge D(dx_{i_j}) \wedge \dots \wedge dx_{i_k} = 0$ b/c $D(dx_{i_j}) \stackrel{①}{=} D(Dx_{i_j}) \stackrel{②}{=} 0$ by hypothesis.

Therefore, if $\omega_p = \sum f_I dx_I$,

$$(D\omega)_p = \sum_I (Df_I \wedge dx_I + f_I \overset{\text{anti-derivation condition}}{D(dx_I)}) \stackrel{①}{=} \sum df_I \wedge dx_I = (d\omega)_p.$$

This is true for all $p \in M$, so $Dw = dw$.

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Ex: If $\omega = f dx + g dy + h dz \in \Omega^1(\mathbb{R}^3)$, then

$$\begin{aligned} d\omega &= \left(\frac{\partial f}{\partial y} dy + \frac{\partial f}{\partial z} dz \right) \wedge dx + \left(\frac{\partial g}{\partial x} dx + \frac{\partial g}{\partial z} dz \right) \wedge dy + \left(\frac{\partial h}{\partial x} dx + \frac{\partial h}{\partial y} dy \right) \wedge dz \\ &= \left(\frac{\partial g}{\partial x} - \frac{\partial f}{\partial y} \right) dx \wedge dy + \left(\frac{\partial f}{\partial z} - \frac{\partial h}{\partial x} \right) dz \wedge dx + \left(\frac{\partial h}{\partial y} - \frac{\partial g}{\partial z} \right) dy \wedge dz. \end{aligned}$$

OTOH, $\omega \longleftrightarrow v = f \frac{\partial}{\partial x} + g \frac{\partial}{\partial y} + h \frac{\partial}{\partial z}$

& $\nabla \times v = \left(\frac{\partial g}{\partial x} - \frac{\partial f}{\partial y} \right) \frac{\partial}{\partial z} + \left(\frac{\partial f}{\partial z} - \frac{\partial h}{\partial x} \right) \frac{\partial}{\partial y} + \left(\frac{\partial h}{\partial y} - \frac{\partial g}{\partial z} \right) \frac{\partial}{\partial x} \longleftrightarrow d\omega$.

So d really does generalize the curl.

Also, if $\eta = a dy \wedge dz + b dz \wedge dx + c dx \wedge dy \longleftrightarrow u = a \frac{\partial}{\partial x} + b \frac{\partial}{\partial y} + c \frac{\partial}{\partial z}$, then

$$d\eta = \frac{\partial a}{\partial x} dx \wedge dy \wedge dz + \frac{\partial b}{\partial y} dy \wedge dz \wedge dx + \frac{\partial c}{\partial z} dz \wedge dx \wedge dy = \left(\frac{\partial a}{\partial x} + \frac{\partial b}{\partial y} + \frac{\partial c}{\partial z} \right) dx \wedge dy \wedge dz \longleftrightarrow \nabla \cdot u = \frac{\partial a}{\partial x} + \frac{\partial b}{\partial y} + \frac{\partial c}{\partial z}.$$

So we've succeeded in defining everything in the following diagram & showed that it commutes (here M is a Riemannian 3-mfld):

$$\begin{array}{ccccccc} \Omega^0(M) & \xrightarrow{d} & \Omega^1(M) & \xrightarrow{d} & \Omega^2(M) & \xrightarrow{d} & \Omega^3(M) \\ \parallel & & \uparrow \iota^* & & \uparrow \iota^* & & \uparrow \iota^* \\ C^0(M) & \xrightarrow{\nabla} & \mathcal{X}(M) & \xrightarrow{\nabla_X} & \mathcal{X}(M) & \xrightarrow{\nabla \cdot} & C^0(M) \end{array}$$

Using the Prop 2(b), if $f: M^n \rightarrow N^n$ is smooth, then we get a map $f^*: \Omega^i(N) \rightarrow \Omega^i(M)$ by

$$(f^* \omega)_p(x_1, \dots, x_k) = \omega_{f(p)}(df_p x_1, \dots, df_p x_k)$$

Prop: $d(f^* \omega) = f^* d\omega$ for any $\omega \in \Omega^k(N)$.

Why is this useful? Well, it means you can compute exterior derivatives in more convenient mds. Later on, we'll consider submersions

$p: SO(n) \rightarrow G_k(\mathbb{R}^n)$ or $g: U(n) \rightarrow G_k(\mathbb{C}^n)$... then we'll be able to compute exterior derivatives y in $SO(n)$ or $U(n)$, where things are nicer.

Proof of Prop: let $p \in M$. If $\omega = u \in \Omega^0(N)$, then $f^*u = u \circ f$, so $d(f^*u) = d(u \circ f)$. On the other hand,

$$(f^*du)_p(X) = du_{f(p)}(df_p X) = (df_p X)(u) = X(u \circ f) = d(u \circ f)_p(X) \text{ for all } X \in T_p M, \text{ so}$$

$f^*du = d(u \circ f)$, completing the proof in the case $k=0$.

Now, suppose $k > 0$ & that $x_1, \dots, x_k$ are local coords in a nbhd of $f(p) \in M$, meaning that $\omega \in \Omega^k(N)$ has the form

$$\omega_{f(p)} = \sum g_I dx_I \text{ near } f(p), \text{ so } (d\omega)_{f(p)} = \sum dg_I \wedge dx_I. \text{ But then}$$

$$(f^*d\omega)_p \stackrel{\text{HW 26a}}{=} \sum (f^*dg_I) \wedge \underset{\text{from } k=0 \text{ case}}{f^*dx_I} = \sum d(g_I \circ f) \wedge f^*dx_{i_1} \wedge \dots \wedge f^*dx_{i_k} \stackrel{\text{from } k=0 \text{ case}}{=} \sum d(g_I \circ f) \wedge d(x_{i_1} \circ f) \wedge \dots \wedge d(x_{i_k} \circ f).$$

OTOH, $(f^*\omega)_p = \sum (g_I \circ f) f^*dx_I = \sum (g_I \circ f) d(x_{i_1} \circ f) \wedge \dots \wedge d(x_{i_k} \circ f)$

so $(df^*\omega)_p = \sum d(g_I \circ f) \wedge d(x_{i_1} \circ f) \wedge \dots \wedge d(x_{i_k} \circ f).$

